## Doping variation of oxygen hole symmetry in layered perovskite nickelates

M. Uchida<sup>1</sup>, Y. Yamasaki<sup>2</sup>, J. Okamoto<sup>2</sup>, Y. Kaneko<sup>3</sup>, H. Nakao<sup>2</sup>, Y. Murakami<sup>2</sup>, and Y. Tokura<sup>1,3,4</sup>

 <sup>1</sup>Department of Applied Physics, University of Tokyo, Tokyo 113-8656, Japan
<sup>2</sup>Photon Factory, Institute of Materials Structure Science, High Energy Accelerator Research Organization, Tsukuba, Ibaraki 305-0801, Japan
<sup>3</sup>Multiferroics Project, ERATO, Japan Science and Technology Agency (JST), Tokyo 113-8656, Japan
<sup>4</sup>Cross-Correlated Materials Research Group (CMRG) and Correlated Electron Research Group (CERG), ASI, RIKEN, Wako 351-0198, Japan

High- $T_{\rm c}$  superconductivity appears close to the Mott transition induced by doping holes into antiferromagnetic parent insulators. Such filling-control insulator-metal transitions are widely observed for transition-metal oxides with strongly correlated electrons, yet the emergence of high-T<sub>c</sub> superconductivity remains unique for the layered cuprates. Layered nickelate  $R_{2-x}Sr_xNiO_4$  (R being rare earth element) with K<sub>2</sub>NiF<sub>4</sub> type structure is a rare example of a two-dimensional antiferromagnetic insulator-metal transition system, providing а contrastive counterpart to superconducting  $R_{2-x}Sr_xCuO_4$  with the same crystal structure. RSNO shows diagonal-stripe and checkerboard charge ordering at  $x \sim 1/3$  and 1/2, respectively, and then undergoes an insulator-metal transition at  $x \sim 1$ .

We have succeeded in growing single crystals of Nd<sub>2-x</sub>Sr<sub>x</sub>NiO<sub>4</sub> up to the metallic region by using a high-pressure floating zone method and investigated the orbital characters of doped holes by systematically measuring polarization-dependent O *K*- and Ni *L*-edge absorption spectra. Figure shows the doping dependence of the O *K*-edge absorption spectra for E || *c*. Two peaks appear above *x*=0.6 and 1.0, respectively, suggesting that the checkerboard type charge ordering persists above *x*=1/2 with introducing the excess holes to  $3z^2 r^2$  orbital states and that the insulator-metal transition occurs with its melting at *x*~1.



FIG. Doping variation of O *K*-edge absorption spectra for E  $\parallel c$  in Nd<sub>2-x</sub>Sr<sub>x</sub>NiO<sub>4</sub>.